

## 2.7 Gravitational waves

### 2.7.1 Weyl tensor

The curvature tensor can be decomposed in terms of the **Weyl tensor** and the Ricci tensor and scalar

$$R_{abcd} = C_{abcd} + \frac{1}{N-2} (R_{ac}g_{bd} - R_{ad}g_{bc} - R_{bc}g_{ad} + R_{bd}g_{ac}) - \frac{1}{(N-1)(N-2)} R (g_{ac}g_{bd} - g_{ad}g_{bc}) \quad (2.7.1)$$

where  $N$  is the dimension of the space. The Ricci tensor and scalar are determined by the matter via the Einstein equation. The Weyl tensor represents gravitational waves.

### 2.7.2 Plane waves

In Rosen coordinates, the metric for a gravitational plane wave is

$$d\tau^2 = 2 du dv - \gamma_{\kappa\lambda}(u) dx^\kappa dx^\lambda \quad (2.7.2)$$

where  $u$  and  $v$  are null coordinates

$$u = \frac{1}{\sqrt{2}}(t - z) \quad (2.7.3)$$

$$v = \frac{1}{\sqrt{2}}(t + z) \quad (2.7.4)$$

and the  $x^\kappa = x, y$  are transverse coordinates. The non-zero Christoffel symbols are

$$\Gamma_{\kappa\lambda}^v = \frac{1}{2}g^{vu}(-g_{\kappa\lambda,u}) = \frac{1}{2}\gamma'_{\kappa\lambda} \quad (2.7.5)$$

$$\Gamma_{u\lambda}^\kappa = \frac{1}{2}g^{\kappa\mu}(g_{\mu\lambda,u}) = \frac{1}{2}\gamma^{\kappa\mu}\gamma'_{\mu\lambda} \quad (2.7.6)$$

plus those related by symmetry. The geodesic equation gives

$$\frac{du}{d\tau} = \text{constant} \quad (2.7.7)$$

$$\frac{d^2v}{d\tau^2} + \frac{1}{2}\gamma'_{\kappa\lambda} \frac{dx^\kappa}{d\tau} \frac{dx^\lambda}{d\tau} = 0 \quad (2.7.8)$$

and

$$\gamma_{\kappa\lambda} \frac{dx^\lambda}{d\tau} = \text{constant} \quad (2.7.9)$$

and so worldlines with  $x^\kappa = \text{constant}$  and  $dz/dt = \text{constant}$  are geodesics despite the fact that the transverse distance between them is varying according to  $\gamma_{\kappa\lambda}(u)$ . The non-zero component of the Einstein tensor is

$$G_{uu} = -\frac{1}{2}\gamma^{\kappa\lambda}\gamma''_{\lambda\kappa} + \frac{1}{4}\gamma^{\kappa\lambda}\gamma'_{\lambda\mu}\gamma^{\mu\nu}\gamma'_{\nu\kappa} \quad (2.7.10)$$

and so the vacuum Einstein equation reduces to

$$\gamma^{\kappa\lambda}\gamma''_{\lambda\kappa} - \frac{1}{2}\gamma^{\kappa\lambda}\gamma'_{\lambda\mu}\gamma^{\mu\nu}\gamma'_{\nu\kappa} = 0 \quad (2.7.11)$$

For small amplitude gravitational waves, we can write

$$\gamma_{\kappa\lambda}(u) = \delta_{\kappa\lambda} + h_{\kappa\lambda}(u) \quad (2.7.12)$$

with  $h_{\kappa\lambda}$  small. Then the Einstein equation linearizes to

$$h''_{xx} + h''_{yy} + O(h^2) = 0 \quad (2.7.13)$$

showing the traceless nature of a gravitational wave.

In Brinkmann coordinates, the metric for a gravitational plane wave is

$$d\tau^2 = 2 du dv + f_{\kappa\lambda}(u) x^\kappa x^\lambda du^2 - \delta_{\kappa\lambda} dx^\kappa dx^\lambda \quad (2.7.14)$$

The non-zero Christoffel symbols are

$$\Gamma_{uu}^\kappa = \frac{1}{2}g^{\kappa\mu}(-g_{uu,\mu}) = \delta^{\kappa\mu}f_{\mu\lambda}x^\lambda \quad (2.7.15)$$

$$\Gamma_{u\kappa}^v = \frac{1}{2}g^{v\mu}(g_{u\mu,\kappa}) = f_{\kappa\lambda}x^\lambda \quad (2.7.16)$$

$$\Gamma_{uu}^v = \frac{1}{2}g^{v\mu}(g_{uu,\mu}) = \frac{1}{2}f'_{\kappa\lambda}x^\kappa x^\lambda \quad (2.7.17)$$

The geodesic equation gives

$$\frac{du}{d\tau} = \text{constant} \quad (2.7.18)$$

$$\frac{d^2v}{d\tau^2} + 2f_{\kappa\lambda}x^\lambda \frac{dx^\kappa}{d\tau} \frac{du}{d\tau} + \frac{1}{2}f'_{\kappa\lambda}x^\kappa x^\lambda \left(\frac{du}{d\tau}\right)^2 = 0 \quad (2.7.19)$$

and

$$\frac{d^2x^\kappa}{d\tau^2} + \delta^{\kappa\mu}f_{\mu\lambda}x^\lambda \left(\frac{du}{d\tau}\right)^2 = 0 \quad (2.7.20)$$

and so worldlines with  $x^\kappa = 0$  and  $dz/dt = \text{constant}$  are geodesics and neighboring geodesics undergo accelerated transverse motion governed by  $f_{\kappa\lambda}(u)$ . The non-zero component of the Einstein tensor is

$$G_{uu} = \delta^{\kappa\lambda}f_{\lambda\kappa} \quad (2.7.21)$$

and so the vacuum Einstein equation reduces to

$$f_{xx} + f_{yy} = 0 \quad (2.7.22)$$

again showing the traceless nature of a gravitational wave.