# Chapter 2

## Electrodynamics

## 2.1 Tensor fields

A **tensor field** is something that takes tensor values at every point in a space, usually relativistic spacetime or Newtonian space and time. Tensor fields of the same type can be added, and multiplied by a scalar, in the usual way. There is also a rich calculus relating different types of tensor fields.

#### 2.1.1 Exterior derivative



Figure 2.1.1: Left: a scalar (zero-form) field gives rise to a covector (one-form) field; right: a covector (one-form) field gives rise to a two-form field.

The relations in Figure 2.1.1 can be mathematically represented as

$$\underline{\nabla} \wedge \phi = \xi \tag{2.1.1}$$

and

$$\underline{\nabla} \wedge \underline{\zeta} = \underline{\underline{\rho}} \tag{2.1.2}$$

where the **exterior derivative**  $\underline{\nabla} \wedge$  has the meaning 'the oriented boundaries of' and gives a measure of the spacial rate of change of the tensor field.

The exterior derivative has the important property

$$\boldsymbol{\nabla} \wedge \boldsymbol{\nabla} \wedge \boldsymbol{\omega} = 0 \tag{2.1.3}$$

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for any differential form  $\omega$ , since the boundary of a boundary is zero, as can be seen from Figure 2.1.1.

### 2.1.2 Integration

An *n*-form field  $\boldsymbol{\omega}$  naturally contracts with an *n*-dimensional surface S to give a scalar

$$\int_{S} \boldsymbol{\omega} = \text{scalar} \tag{2.1.4}$$

with the same interpretation as the contraction of an *n*-form with an *n*-vector. If we divide the surface S into infinitesimal surface elements dS, the integral of  $\omega$  over S can be written in the more familiar form

$$\int_{S} \boldsymbol{\omega} \cdot \boldsymbol{dS} \tag{2.1.5}$$

#### 2.1.3 Stokes' theorem

Stokes' theorem states that

$$\int_{S} \boldsymbol{\nabla} \wedge \boldsymbol{\omega} = \int_{\partial S} \boldsymbol{\omega}$$
(2.1.6)

where  $\partial S$  is the boundary of S.



Figure 2.1.2: Stokes' theorem:  $\int_{S} \nabla \wedge \boldsymbol{\omega} = \int_{\partial S} \boldsymbol{\omega} = 2.$