

Homework 9

Q9.1. Consider the differential operator

$$L_x = \frac{d^2}{dx^2} \quad (\text{Q9.1.1})$$

acting on functions $\phi : [0, 2\pi] \rightarrow \mathbb{C}$.

What is the most general boundary condition consistent with L_x being Hermitian? Show that

i. the Dirichlet boundary conditions

$$\phi(0) = \phi(2\pi) = 0 \quad (\text{Q9.1.2})$$

ii. the Neumann boundary conditions

$$\phi'(0) = \phi'(2\pi) = 0 \quad (\text{Q9.1.3})$$

iii. the periodic boundary conditions

$$\phi(0) = \phi(2\pi) \quad \text{and} \quad \phi'(0) = \phi'(2\pi) \quad (\text{Q9.1.4})$$

are all special cases of this general boundary condition.

In each case

- determine the eigenvalues and eigenspaces,
- state the orthogonality and completeness equations,
- check that $W(\phi, \psi) = \text{constant}$ for ϕ and ψ in the same eigenspace,
- express

$$f(x) = 2\pi x - x^2 \quad (\text{Q9.1.5})$$

in terms of eigenvectors of L_x .

A9.1.

$$\phi^* L_x \psi - \psi L_x^* \phi^* = \frac{d}{dx} \left(\phi^* \frac{d\psi}{dx} - \psi \frac{d\phi^*}{dx} \right) \quad (\text{A9.1.1})$$

Therefore, from Eqs. (3.4.8) to (3.4.10), the general Hermitian boundary condition is

$$\phi^*(0) \psi'(0) - \psi(0) \phi^{*'}(0) = \phi^*(2\pi) \psi'(2\pi) - \psi(2\pi) \phi^{*'}(2\pi) \quad (\text{A9.1.2})$$

The Dirichlet, Neumann and periodic boundary conditions, Eqs. (Q9.1.2), (Q9.1.3) and (Q9.1.4), are clearly special cases of this general boundary condition.

(a) The general solution of the eigenvector equation

$$L_x \phi = \lambda \phi \quad (\text{A9.1.3})$$

is

$$\phi_\lambda(x) = \begin{cases} A_0 + B_0 x & \text{for } \lambda = 0 \\ A_\lambda e^{\sqrt{\lambda}x} + B_\lambda e^{-\sqrt{\lambda}x} & \text{for } \lambda \neq 0 \end{cases} \quad (\text{A9.1.4})$$

i. The Dirichlet boundary conditions Eq. (Q9.1.2) constrain the coefficients to

$$A_0 = B_0 = 0 \quad (\text{A9.1.5})$$

and

$$A_\lambda + B_\lambda = 0 \quad (\lambda \neq 0) \quad (\text{A9.1.6})$$

and the eigenvalues to

$$\lambda = -\frac{n^2}{4} \quad (n \in \mathbb{Z}, n > 0) \quad (\text{A9.1.7})$$

Therefore the eigenspaces are

$$\phi_\lambda(x) = C_\lambda \sin\left(\frac{nx}{2}\right) \quad (\text{A9.1.8})$$

ii. The Neumann boundary conditions Eq. (Q9.1.3) constrain the coefficients to

$$B_0 = 0 \quad (\text{A9.1.9})$$

and

$$A_\lambda - B_\lambda = 0 \quad (\lambda \neq 0) \quad (\text{A9.1.10})$$

and the eigenvalues to

$$\lambda = -\frac{n^2}{4} \quad (n \in \mathbb{Z}, n \geq 0) \quad (\text{A9.1.11})$$

Therefore the eigenspaces are

$$\phi_\lambda(x) = C_\lambda \cos\left(\frac{nx}{2}\right) \quad (\text{A9.1.12})$$

iii. The periodic boundary conditions Eq. (Q9.1.4) constrain the coefficients to

$$B_0 = 0 \quad (\text{A9.1.13})$$

and the eigenvalues to

$$\lambda = -n^2 \quad (n \in \mathbb{Z}) \quad (\text{A9.1.14})$$

Therefore the eigenspaces are

$$\phi_\lambda(x) = A_\lambda e^{inx} + B_\lambda e^{-inx} \quad (\text{A9.1.15})$$

(b) i. Orthogonality

$$\int_0^{2\pi} dx \sin\left(\frac{mx}{2}\right) \sin\left(\frac{nx}{2}\right) = \pi \delta_{mn} \quad (\text{A9.1.16})$$

Completeness

$$\sum_{n=1}^{\infty} \sin\left(\frac{nx}{2}\right) \sin\left(\frac{ny}{2}\right) = \pi [\delta_{4\pi}(x-y) - \delta_{4\pi}(x+y)] \quad (\text{A9.1.17})$$

where

$$\delta_p(x) = \delta(x) \quad \left(-\frac{p}{2} \leq x \leq \frac{p}{2}\right) \quad (\text{A9.1.18})$$

and

$$\delta_p(x) = \delta_p(x+p) \quad (\text{A9.1.19})$$

ii. Orthogonality

$$\int_0^{2\pi} dx \cos\left(\frac{mx}{2}\right) \cos\left(\frac{nx}{2}\right) = \pi \delta_{mn} + \pi \delta_{m0} \delta_{n0} \quad (\text{A9.1.20})$$

Completeness

$$\frac{1}{2} + \sum_{n=1}^{\infty} \cos\left(\frac{nx}{2}\right) \cos\left(\frac{ny}{2}\right) = \pi [\delta_{4\pi}(x-y) + \delta_{4\pi}(x+y)] \quad (\text{A9.1.21})$$

iii. Orthogonality

$$\int_0^{2\pi} dx e^{-imx} e^{inx} = 2\pi \delta_{mn} \quad (\text{A9.1.22})$$

Completeness

$$\sum_{n=-\infty}^{\infty} e^{inx} e^{-iny} = 2\pi \delta_{2\pi}(x-y) \quad (\text{A9.1.23})$$

(c)

$$W(\phi, \psi) = \phi^* \psi' - \phi'^* \psi \quad (\text{A9.1.24})$$

i.

$$W\left(\sin \frac{nx}{2}, \sin \frac{nx}{2}\right) = 0 \quad (\text{A9.1.25})$$

ii.

$$W\left(\cos \frac{nx}{2}, \cos \frac{nx}{2}\right) = 0 \quad (\text{A9.1.26})$$

iii.

$$W(e^{inx}, e^{inx}) = 2in \quad (\text{A9.1.27})$$

$$W(e^{-inx}, e^{inx}) = 0 \quad (\text{A9.1.28})$$

(d) i.

$$f(x) = \sum_{n=1}^{\infty} f_n \sin\left(\frac{nx}{2}\right) \quad (\text{A9.1.29})$$

where

$$f_n = \frac{1}{\pi} \int_0^{2\pi} dx \sin\left(\frac{nx}{2}\right) f(x) \quad (\text{A9.1.30})$$

$$= \begin{cases} \frac{32}{\pi n^3} & \text{for } n \text{ odd} \\ 0 & \text{for } n \text{ even} \end{cases} \quad (\text{A9.1.31})$$

ii.

$$f(x) = \frac{f_0}{2} + \sum_{n=1}^{\infty} f_n \cos\left(\frac{nx}{2}\right) \quad (\text{A9.1.32})$$

where

$$f_n = \frac{1}{\pi} \int_0^{2\pi} dx \cos\left(\frac{nx}{2}\right) f(x) \quad (\text{A9.1.33})$$

$$= \begin{cases} \frac{4\pi^2}{3} & \text{for } n = 0 \\ 0 & \text{for } n \text{ odd} \\ -\frac{16}{n^2} & \text{for } n \text{ even, } n > 0 \end{cases} \quad (\text{A9.1.34})$$

iii.

$$f(x) = \sum_{n=-\infty}^{\infty} f_n e^{inx} \quad (\text{A9.1.35})$$

where

$$f_n = \frac{1}{2\pi} \int_0^{2\pi} dx e^{-inx} f(x) \quad (\text{A9.1.36})$$

$$= \begin{cases} \frac{2\pi^2}{3} & \text{for } n = 0 \\ -\frac{2}{n^2} & \text{for } n \neq 0 \end{cases} \quad (\text{A9.1.37})$$

Q9.2. The real differential equation

$$\frac{d^2 y}{dx^2} + a(x) \frac{dy}{dx} + [\lambda b(x) + c(x)] y = 0 \quad (\text{Q9.2.1})$$

with boundary conditions

$$y(0) = y(1) = 0 \quad (\text{Q9.2.2})$$

has solutions $y_\lambda(x)$. What properties do the $y_\lambda(x)$ have?

A9.2. Rearranging Eq. (Q9.2.1) gives the eigen equation

$$L_x y = \lambda y \quad (\text{A9.2.1})$$

with

$$L_x = -\frac{1}{b(x)} \frac{d^2}{dx^2} - \frac{a(x)}{b(x)} \frac{d}{dx} - \frac{c(x)}{b(x)} \quad (\text{A9.2.2})$$

Assuming $a(x)$, $b(x)$ and $c(x)$ are real, and using Eqs. (3.4.18) to (3.4.21), L_x has the form of a Hermitian differential operator

$$g(x) L_x = -\frac{d}{dx} p(x) \frac{d}{dx} - q(x) \quad (\text{A9.2.3})$$

with

$$p(x) = e^{\int dx a(x)} \quad (\text{A9.2.4})$$

$$q(x) = c(x) e^{\int dx a(x)} \quad (\text{A9.2.5})$$

$$g(x) = b(x) e^{\int dx a(x)} \quad (\text{A9.2.6})$$

Therefore the eigenfunctions $y_\lambda(x)$ will be orthogonal¹, over the range determined by the boundary conditions Eq. (Q9.2.2) and with the integration measure of Eq. (A9.2.6),

$$\int_0^1 dx b(x) e^{\int dx a(x)} y_\lambda^*(x) y_{\lambda'}(x) \propto \delta_{\lambda\lambda'} \quad (\text{A9.2.7})$$

and complete

$$\sum_\lambda \frac{y_\lambda(x) y_\lambda^*(x')}{\int_0^1 dx'' b(x'') e^{\int dx'' a(x'')} y_\lambda^*(x'') y_\lambda(x'')} = \frac{\delta(x, x')}{b(x) e^{\int dx a(x)}} \quad (\text{A9.2.8})$$

¹Degenerate eigenfunctions can be chosen to be orthogonal, and will also satisfy Eq. (3.4.11). However, the notation $y_\lambda(x)$ would seem to imply that they are not degenerate.